AWFM06/0053 Bifurcation in Resistive Drift Wave Turbulence <u>Ryusuke Numata*</u>, Rowena Ball, Robert L. Dewar, Linda Stals^{A)}, Claudio Tebaldi^{B)}

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In this work we report the results of numerical simulation of the modified Hasegawa-Wakatani model, which describes resistive drift wave turbulence in tokamak edge plasmas. The HW model is a simple model, but contains an interesting physics, i.e. turbulence-shear flow interaction. This simulation study also complements the low-dimensional modeling results which we already have.

Modified Hasegawa-Wakatani Model

Stability Diagram Indicating 1st Bifurcation Points

MHW model describes evolution of density fluctuation n and vorticity $\zeta = \nabla^2 \varphi$ (φ : electrostatic potential)

 $\frac{\partial \zeta}{\partial t} + \{\varphi, \zeta\} = \alpha (\tilde{\varphi} - \tilde{n}) - D_{\zeta} \nabla^{4} \zeta$ $\frac{\partial n}{\partial t} + \{\varphi, n\} = \alpha (\tilde{\varphi} - \tilde{n}) - \kappa \frac{\partial \varphi}{\partial y} - D_{n} \nabla^{4} n$

 $[a, b] = \partial a / \partial x \partial b / \partial y - \partial a / \partial y \partial b / \partial x, \nabla^2 = \partial^2 / \partial x^2 + \partial^2 / \partial y^2$ D_{τ} , D_n : dissipation coefficients (Prandtl number $P_r = D_{\tau}/D_n$)

$$\kappa = \frac{-\partial}{\partial x} \ln n_0, \quad \alpha = \frac{T_e k_z^2}{\eta n_0 \omega_{ci} e^2}$$

Resistive coupling term which comes from parallel electron response $\partial j_z / \partial z = 1/\eta \partial^2 (\varphi - n) / \partial z^2$ (Ohm's law) does not act on the zonal components $(k_v = k_z = 0)$

> Original HW MHW $\alpha(\varphi - n) \rightarrow \alpha(\tilde{\varphi} - \tilde{n})$

Zonal:
$$\langle f \rangle = \frac{1}{L_y} \int f \, dy$$
 (*f* stands for φ and *n*)
Non-zonal: $\tilde{\varphi} = \varphi - \langle \varphi \rangle$, $\tilde{n} = n - \langle n \rangle$





For finite α one positive growth rate if D=0 (dissipationless) large $\kappa \rightarrow$ destabilizing, large $\alpha \rightarrow$ stabilizing Growth rate monotonically decrease with k_x , peaks at $k_v \sim 1$

Numerical Algorithm to Solve MHW Model

•MHW model is solved in 2D slab geometry •Box size L, determined by smallest wavenumber $\Delta k=0.3 [(2L)^2=(2\pi/\Delta k)^2]$ •Periodic boundary in both x and y direction •Time stepping algorithm is a 3rd order explicit linear multistep method •Finite difference method is used for spatial discretization •Poisson bracket term evaluated by the Arakawa's method (Arakawa(1966)) •Implemented on APAC-NF SGI Altix 3700 Bx2 Cluster

Zonal Flow Generation and Transport Suppression

Transition to Zonal Flow Suppressed Regime

Contour plot of saturated variables for $\kappa = 1$, $\alpha = 1$, $D_{\tau} = 10^{-6}$, $P_r = 1$











To understand Bifurcation ...

- •Stability of Zonal Flow [Kelvin-Helmholtz Instability]
 - $\varphi = \varphi_0(x) + \tilde{\varphi}(x) \exp(ik_v y + \gamma t)$

•What is the 1st destabilized structure which appears just beyond linear threshold?

Fully nonlinear





Summary and Conclusion

- •We have performed simulations of modified Hasegawa-Wakatani model •Zonal flows are generated, and the zonal flows suppress cross-field turbulent transport in MHW model
- •Transition from zonal flow dominant to zonal flow suppressed state is observed •Bifurcation points may be understand as destabilization of zonal flow or some other structure. But, what structure?
- •To look for the 1st destabilized state, we start to study the behavior near the linear stability boundary where only few modes are involved.
- •Oscillatory behavior is observed just beyond linear stability boundary, where drift wave and zonal flow are competing [predator-pray]



•Corresponence of bifurcation diagram obtained from numerical simulation and from low- *k* dimensional model should be considered

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