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Introduction

Magnetic reconnection is the phenomena which couples widely separated scales
 Using gyrokinetics, we can seamlessly study different regimes in a strong guide field limit

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Derivation of Ohm's law

Gyrokinetic equation

describes evolutions of distribution function h in gyrocenter coordinate (\mathbf{R}, \mathbf{V}) $\mathbf{R} = \mathbf{r} + \frac{\mathbf{v} \times \mathbf{b}_0}{\Omega}$

$$f_s = f_{0s} + q_s \phi T_{0s} + h(\mathbf{R})$$

$$\frac{\partial h_s}{\partial t} + \mathbf{V} \cdot \frac{\partial h_s}{\partial \mathbf{Z}} + \frac{1}{B_0} \langle \mathbf{X} \rangle_{\mathbf{R}, s} \cdot \nabla h_s = \frac{q_s f_{0s}}{T_{0s}} \frac{\partial \langle \mathbf{X} \rangle_{\mathbf{R}}}{\partial t} + \langle C(h_s) \rangle_{\mathbf{R}, s}$$

$$\mathbf{X} = \phi - \mathbf{v} \cdot \mathbf{A}$$

Gyrokinetic field equations

relates between fields and moments in \mathbf{r}

$$\sum_s q_s \tilde{n}_s = 0, \quad \tilde{n}_s = \int \delta f_s d\mathbf{v} = \int \langle h_s \rangle_r d\mathbf{v} - \frac{q_s n_{0s}}{T_{0s}} \phi$$

$$\nabla_{\perp}^2 A_{\parallel} = -\mu_0 \sum_s q_s n_{0s} \tilde{u}_{\parallel, s}, \quad \tilde{u}_{\parallel, s} = \frac{1}{n_{0s}} \int \delta f_s v_{\parallel} d\mathbf{v} = \frac{1}{n_{0s}} \int \langle h_s v_{\parallel} \rangle_r d\mathbf{v}$$

$$B_0 \nabla_{\perp} \delta B_{\parallel} = -\mu_0 \nabla_{\perp} \cdot \tilde{\mathbf{P}}_{\perp, s}, \quad \tilde{\mathbf{P}}_{\perp, s} = \int m \mathbf{v} \mathbf{v} \delta f_s d\mathbf{v}$$

$$= \int m v_{\parallel} \langle h_s \rangle_r d\mathbf{v} + \int \langle m v_{\perp} v_{\perp} h_s \rangle_r d\mathbf{v} - q_s n_{0s} \phi \mathbf{I}$$

$$= P_{\parallel, s} \mathbf{e}_{\parallel} \mathbf{e}_{\parallel} + P_{\perp, s} \mathbf{e}_{\perp} \mathbf{e}_{\perp} + P_{\theta, s} \mathbf{e}_{\theta} \mathbf{e}_{\theta} + P_{\phi, s} \mathbf{e}_{\phi} \mathbf{e}_{\phi}$$

Ohm's law

is derived by taking 1st order v_{\parallel} moment in \mathbf{r}

$$m_e n_{0e} \frac{\partial \tilde{u}_{\parallel, e}}{\partial t} + \frac{1}{B_0} (m_e n_{0e} G(\langle \phi, \tilde{u}_{\parallel, e} \rangle) - G(\langle A_{\parallel}, (\tilde{\mathbf{P}}_{\parallel, e} + n_{0e} q_e \phi) \rangle) + G(\nabla^{-1} \langle \delta B_{\parallel}, \tilde{\mathbf{P}}_{\theta, e} \rangle)) + n_{0e} q_e \frac{\partial \langle \langle A_{\parallel} \rangle_r \rangle}{\partial t} = M_{\parallel}$$

Inertia Off-diagonal pressure

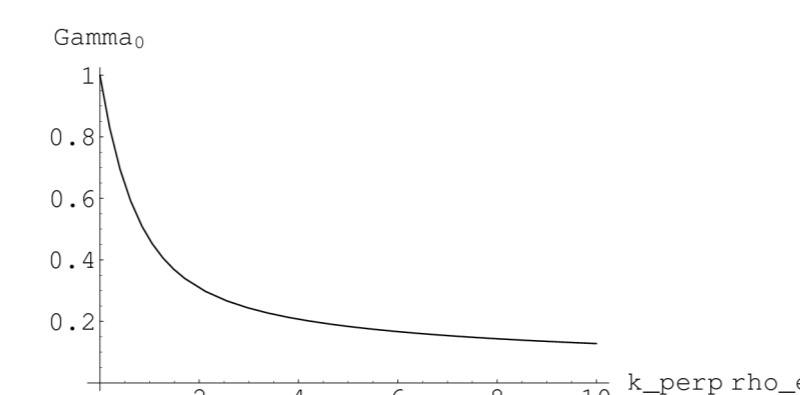
Where nonlinear terms cannot be simply given by fields and moments

$$G(\langle \phi, \tilde{u}_{\parallel, e} \rangle) = \sum_k e^{i\mathbf{k} \cdot \mathbf{r}} \sum_{\mathbf{k}'' = -\mathbf{k}} \left(-[\mathbf{k}', \mathbf{k}''] \phi_{\mathbf{k}'} \frac{1}{n_0} \int J_0(\alpha_e' + \alpha_e'') J_0(\alpha_e') v_{\parallel} h_{e, \mathbf{k}''} d\mathbf{v} \right)$$

$$G(\langle A_{\parallel}, (\tilde{\mathbf{P}}_{\parallel, e} + n_{0e} q_e \phi) \rangle) = \sum_k e^{i\mathbf{k} \cdot \mathbf{r}} \sum_{\mathbf{k}'' = -\mathbf{k}} \left(-[\mathbf{k}', \mathbf{k}''] A_{\parallel, \mathbf{k}'} \int J_0(\alpha_e' + \alpha_e'') J_0(\alpha_e') m v_{\parallel}^2 h_{e, \mathbf{k}''} d\mathbf{v} \right)$$

$$G(\nabla^{-1} \langle \delta B_{\parallel}, \tilde{\mathbf{P}}_{\theta, e} \rangle) = \sum_k e^{i\mathbf{k} \cdot \mathbf{r}} \sum_{\mathbf{k}'' = -\mathbf{k}} \left(-[\mathbf{k}', \mathbf{k}''] \frac{\delta B_{\parallel, \mathbf{k}'}}{i k_{\perp}'} \int i J_1(\alpha_e' + \alpha_e'') J_0(\alpha_e') m v_{\perp} v_{\parallel} h_{e, \mathbf{k}''} d\mathbf{v} \right)$$

Electron feels $\langle \langle A_{\parallel} \rangle_r \rangle = \sum_k e^{i\mathbf{k} \cdot \mathbf{r}} A_{\parallel, \mathbf{k}} \Gamma_0 \left(\frac{k_{\perp}^2 \rho_c^2}{2} \right)$



Limiting cases

In $k_{\perp} \rho_c \ll 1$ limit¹⁾,

$$m_e n_{0e} \frac{\partial \tilde{u}_{\parallel, e}}{\partial t} + \frac{1}{B_0} (m_e n_{0e} \langle \phi, \tilde{u}_{\parallel, e} \rangle - [A_{\parallel}, (\tilde{\mathbf{P}}_{\parallel, e} + n_{0e} q_e \phi)] + \langle \nabla^{-1} \delta B_{\parallel}, \tilde{\mathbf{P}}_{\theta, e} \rangle) + n_{0e} q_e \frac{\partial \langle A_{\parallel} \rangle_r}{\partial t} = M_{\parallel}$$

For perturbed maxwellian distribution

$$h_e = f_{0e} \left(\frac{\tilde{n}_e}{n_{0e}} + \frac{q_e \phi}{T_{0e}} + \left(\frac{v^2}{v_{th, e}^2} - \frac{3}{2} \right) \frac{\tilde{T}_e}{T_{0e}} + \frac{2 v_{\parallel}}{v_{th, e}} \frac{\tilde{u}_{\parallel}}{v_{th, e}} \right)$$

$$P_{\parallel} = P_{\perp} = P_{\theta} = \tilde{n}_e T_{0e} + n_{0e} \tilde{T}_e, P_{\phi} = 0$$

Using the perpendicular Ampere's law, we recover the Hall term from

AstroGK²⁾

(<http://www.physics.uiowa.edu/~ghowes/astrogk/>)

We are developing a slab gyrokinetics code AstroGK (a simplified version of GS2). Recently, we have implemented full electron-ion collision, which was missing in the previous work using GS2³⁾

$$\langle C_{ei}(h) \rangle_r = \sum_k e^{i\mathbf{k} \cdot \mathbf{r}} v_{ei} \left(\frac{v_{th, e}}{V} \right)^3 \left(\frac{1}{2} \frac{\partial}{\partial \xi} (1 - \xi^2) \frac{\partial h_{e, \mathbf{k}}}{\partial \xi} - \frac{k^2 V^2}{4 \Omega_{0e}} (1 + \xi^2) h_{e, \mathbf{k}} + \frac{2 J_0(\alpha_e) V_{\parallel} u_{\parallel, i}}{v_{th, e}^2} f_{0e} \right)$$

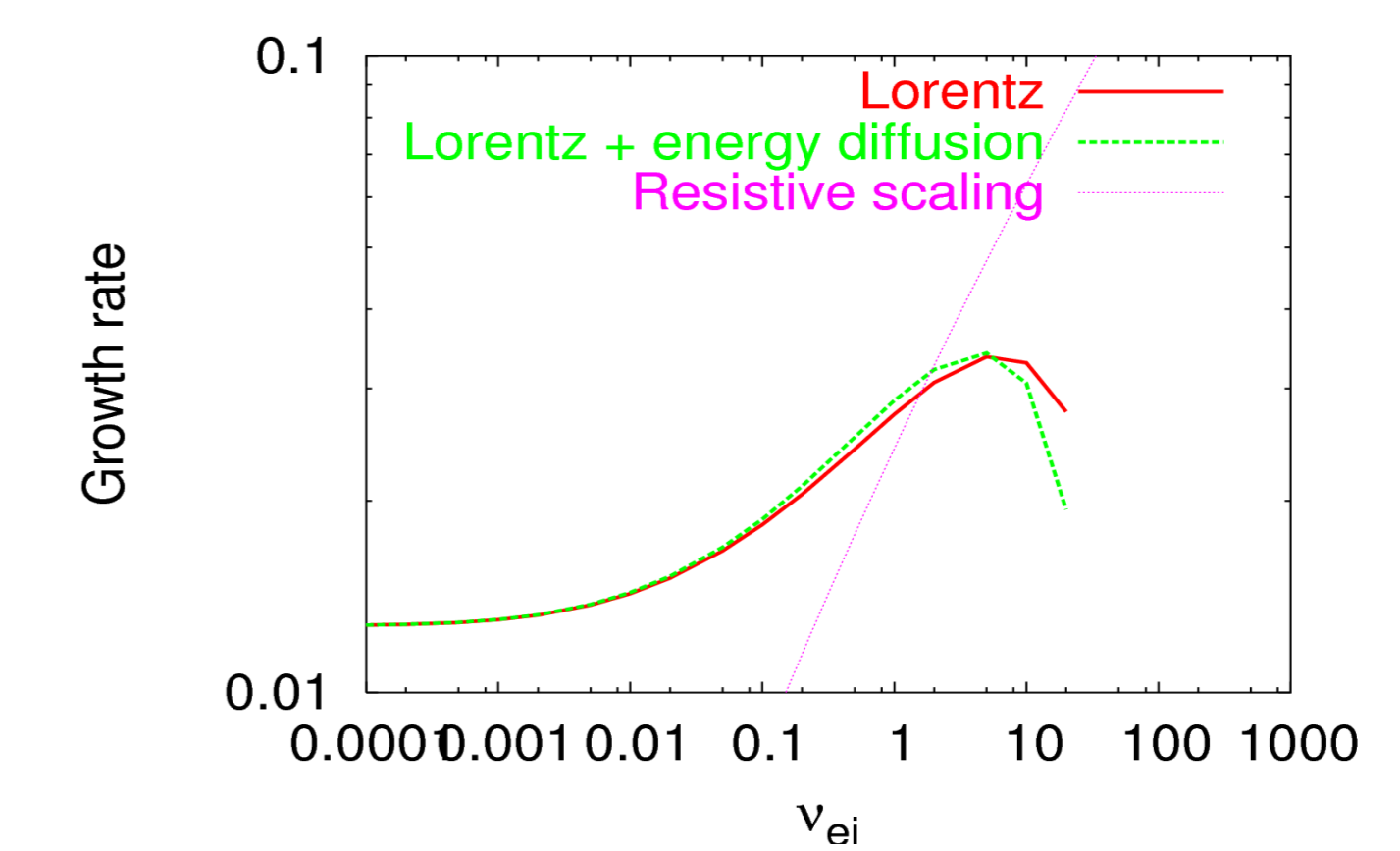
Numerical test of the e-i collision term using the classical Spitzer problem shows that the resistivity is given by

$$\eta = m_e / (1.98 \tau_e n_e e^2)$$

where $\tau_e = 3\sqrt{\pi}/4 v_{ei}$ in $k_{\perp} \rho_c \ll 1$ limit⁴⁾

Tearing mode

Collisionless regime

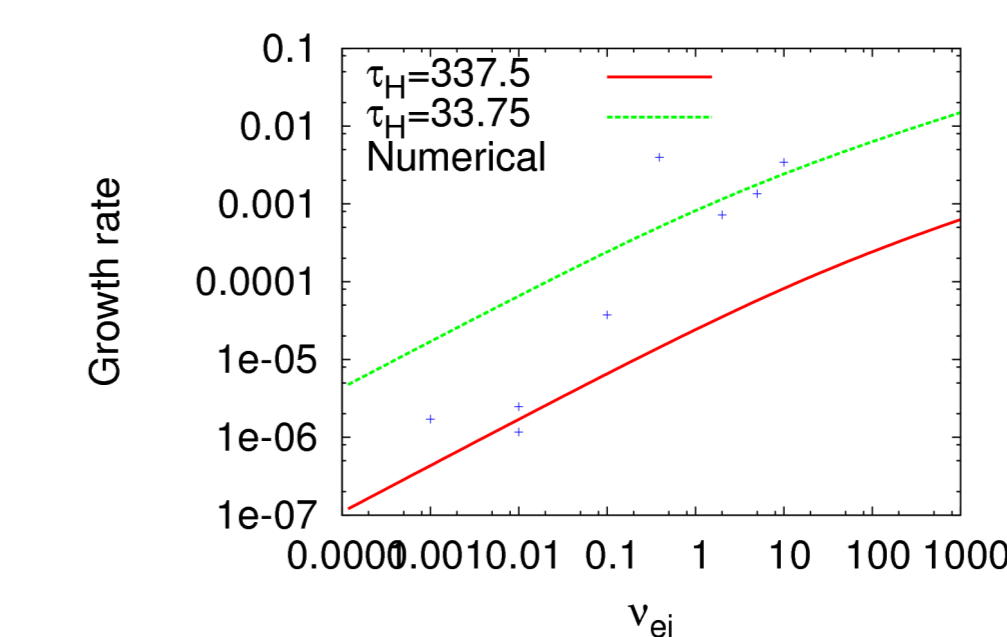


$$\Delta' = 23$$

$$\tau_H = 3.375$$

$$\tau_R = \frac{1965}{v_{ei}}$$

Collisional regime



Conclusions

We have derived ohm's law for gyrokinetic electrons. Using recently implemented electron-ion collision in AstroGK, we have performed tearing mode simulations for collisionless and collisional regimes.

References

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